

# On the energy estimates for second order homogeneous hyperbolic equations with Levi-type conditions <sup>1</sup>

Fumihiko Hirosawa (Yamaguchi University)

and

Bui Tang Bao Ngoc (Hanoi University of Technology)

We consider the following Cauchy problem of a second order homogeneous hyperbolic equation with variable coefficients:

$$\begin{cases} (\partial_t^2 - a(t)^2 \partial_x^2 + 2b(t) \partial_x \partial_t) u = 0, & (t, x) \in (0, \infty) \times \mathbb{R}, \\ (u(0, x), \partial_t u(0, x)) = (u_0(x), u_1(x)), & x \in \mathbb{R}, \end{cases} \quad (1)$$

where  $a(t), b(t) \in C^m([0, \infty))$  ( $m \geq 2$ ) are real valued and satisfy the following strictly hyperbolic condition:

$$0 < c_0 \leq c(t) := \sqrt{a(t)^2 + b(t)^2} \leq c_1. \quad (2)$$

Here we introduce the following energy to the solution of (1):

$$E(t) := \frac{1}{4} \int_{\mathbb{R}} \left( |\partial_t u(t, x) + (b(t) + c(t)) \partial_x u(t, x)|^2 + |\partial_t u(t, x) + (b(t) - c(t)) \partial_x u(t, x)|^2 \right) dx. \quad (3)$$

If the coefficients are constants, then the energy conservation  $E(t) \equiv E(0)$  is valid. However, we cannot expect such a property for variable coefficients; thus we introduce the following property of an equivalence of the energy with respect to  $t$ :

$$C^{-1}E(0) \leq E(t) \leq CE(0), \quad (\text{GEC})$$

which is called the *generalized energy conservation*, where  $C > 1$  is a constant.

If  $b(t) = 0$ , then the equation of (1) is a wave equation with a variable propagation speed, and  $a'(t)$  describes the oscillating speed of it. Trivially, we see that (GEC) is valid if  $a' \in L^1(\mathbb{R}_+)$  though  $b(t) \neq 0$ . However, it is not clear whether (GEC) holds or not if  $a' \notin L^1(\mathbb{R}_+)$ . Actually, (GEC) is not true in general; indeed, for  $b(t) = 0$  an example of  $a(t)$  is constructed in [6]. The main purpose of our research is to have some conditions to the coefficients which provide (GEC) taking account of the  $C^m$  regularity of the coefficients. In particular, we focus the conditions between  $a(t)$  and  $b(t)$ , which give the same conclusion of [3] considering (1) with  $b(t) = 0$ .

Let us introduce the following conditions to the coefficients:

- *Stabilization condition*: there exist the means  $a_\infty$  and  $b_\infty$  of  $a(t)$  and  $b(t)$  on  $\mathbb{R}_+$  such that

$$\int_0^t (|a(s) - a_\infty| + |b(s) - b_\infty|) ds \leq C_0(1+t)^\alpha \quad \text{for } \alpha \in [0, 1]. \quad (4)$$

- *Control of the oscillations*:

$$|a^{(k)}(t)| + |b^{(k)}(t)| \leq C a k e (1+t)^{-k\beta} \quad \text{for } \beta \in [0, 1] \quad (k = 1, \dots, m). \quad (5)$$

REMARK 1. The stabilization condition (4) is trivial for  $\alpha = 1$  since (2) is valid. The condition of control of the oscillations (5) with  $\beta > 1$  with  $k = 1$  gives  $a'(t) \in L^1(\mathbb{R}_+)$ .

<sup>1</sup>The 26th Matsuyama camp, "Dedicated to the 60th Birthday of Professor Morimoto", January 6-9, 2011, Kyoto.

Let us recall the following result in [3] for  $b(t) = 0$ :

**Theorem 1** ([3]). *Let  $b(t) = 0$  and  $m \geq 2$ . If  $a(t)$  satisfies (2), (4) and (5) for*

$$\beta = \beta_m := \alpha + \frac{1 - \alpha}{m}, \quad (6)$$

*then (GEC) is valid.*

REMARK 2. The restriction to the order of  $a'(t)$  in (5) is weaker as  $\beta$  larger, which is realized as  $m$  larger. That is, faster oscillation to the coefficient is possible to be permissible for (GEC) as the coefficient is smoother. Here we underline that we have a benefit by the choice of larger  $m$  only for  $\alpha < 1$ ; thus the stabilization property (4) is essential.

It may be natural that we expect the same conclusion of Theorem 1 for  $b(t) \neq 0$ . However, we see from the analogy of the result in [4] that such an expectation is not valid to the general model (1) with  $b(t) \neq 0$ , because an interaction between the oscillating coefficients  $a(t)$  and  $b(t)$  gives a bad effect for (GEC). On the other hand, the result in [5] hints us that the following condition between  $a(t)$  and  $b(t)$ :

$$\sup_t \left\{ \left| \int_0^t \frac{b'(s)}{c(s)} ds \right| \right\} \leq C, \quad (L1)$$

which is called the  $C^1$ -type Levi condition, is possible to invalidate the bad effect from the interactions of the oscillating coefficients. Indeed, we have the following theorem:

**Theorem 2** ( $C^3$  coefficients [1]). *Let  $m = 2$  or  $3$ ,  $\alpha \in [0, 1)$  and  $\beta = \beta_m$ . Assume that  $a, b \in C^m([0, \infty))$  satisfy (2), (4) and (5). If the  $C^1$ -type Levi condition (L1) holds, then the generalized energy conservation (GEC) is valid.*

REMARK 3. (L1) is true if  $a(t)$  is represented by  $a(t) = \phi(b(t))$  with a positive  $C^1$  function  $\phi$ .

REMARK 4. Actually, under the assumption (L1) one can prove (GEC) for  $m = 2$  and  $\beta (= \beta_b) = 1$  (see [2, 7], which consider more general hyperbolic systems and  $L^p$ - $L^q$  type decay estimates).

We cannot have the same conclusion as Theorem 2 for  $m \geq 4$ . However, we have the following theorem for (GEC) with  $m = 4, 5$  if we additionally suppose the  $C^2$ -type Levi condition:

**Theorem 3** ( $C^5$  coefficients [1]). *Let  $m = 4$  or  $5$ ,  $\alpha \in [0, 1)$  and  $\beta = \beta_m$ , where  $\beta_m$  is defined by (6). Assume that  $a, b \in C^m([0, \infty))$  satisfy (2), (4) and (5). If the  $C^1$ -type Levi condition (L1) and the  $C^2$ -type Levi condition:*

$$\sup_t \left\{ (1+t)^{2\alpha} \left| \int_0^t \frac{c(s) (b'(s)c''(s) - b''(s)c'(s)) - b'(s) ((b'(s))^2 - (c'(s))^2)}{c(s)^5} ds \right| \right\} \leq C \quad (L2)$$

*hold, then the generalized energy conservation (GEC) is valid.*

It is a natural observation that Theorem 2 and 3 may be generalized for  $m = 6, 7$ ,  $m = 8, 9$ , and so on under some  $C^k$ -type Levi conditions with  $k = 3, 4, \dots$ . Actually, it is true in a certain sense, but the representations of the corresponding  $C^k$ -type Levi conditions are very complicate.

## References

- [1] T. B. N. Bui, F. Hirosawa, On the energy estimates for second order hyperbolic equations with time dependent coefficients. Preprint.
- [2] M. D’Abbicco, S. Luccente, G. Taglialatela,  $L^p$ - $L^q$  estimates for regularly hyperbolic systems. *Adv. Differential Equations* **14** (2009), 801–834.
- [3] F. Hirosawa, On the asymptotic behavior of the energy for the wave equations with time depending coefficients. *Math. Ann.* **339** (2007), 819–839.
- [4] F. Hirosawa, M. Reissig, About the optimality of oscillations in non-Lipschitz coefficients for strictly hyperbolic equations. *Ann. Scuola Norm. Sup. Pisa Cl. Sci. (5)* **3** (2004), 589–608.
- [5] F. Hirosawa, M. Reissig, Levi condition for hyperbolic equations with oscillating coefficients. *J. Differential Equations.* **223** (2006), 329–350.
- [6] M. Reissig, J. Smith,  $L^p$ - $L^q$  estimate for wave equation with bounded time dependent coefficient. *Hokkaido Math. J.* **34** (2005), 541–586.
- [7] M. Ruzhansky, J. Wirth, Corrigendum to “dispersive estimates for  $T$ -dependent hyperbolic systems”. *Rend. Sem. Mat. Univ. Politec. Torino* **68** (2010), 339–349.