Energy estimates for wave equations with time dependent coefficients

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1. Introduction

Consider the following Cauchy problem and the total energy:

$$\begin{cases} (\partial_t^2 - a(t)^2 \Delta) \ u(t, x) = 0, \ (t, x) \in [0, \infty) \times \mathbb{R}^n, \\ (u(0, x), u_t(0, x)) = (u_0(x), u_1(x)), \ x \in \mathbb{R}^n, \end{cases}$$
(1)

 $0 < a_0 \le a(t) \le a_1, a(t) \in C^1([0,\infty)).$

$$E(t) = \frac{1}{2} \left(|a(t)|^2 ||\nabla u(t, \cdot)||^2 + ||\partial_t u(t, \cdot)||^2 \right)$$
 (2)

If a(t) is a constant, then the energy conservation (EC) is valid:

$$E(t) \equiv E(0),$$
 (EC)

however, (EC) is not valid if a(t) is not a constant.

Generally, we only expect the following estimates:

$$\eta(t)^{-1}E(0) \le E(t) \le \eta(t)E(0) \quad (\eta(t) > 1).$$
 (3)

 $\eta(t) \equiv C \Leftrightarrow \text{GEC}(=\text{Generalized Energy Conservation}).$

Monotone increasing or decreasing coefficients

$$E'(t) = a'(t)a(t)\|\nabla u(t)\|^2$$

$$a'(t) \leq 0 \ \Rightarrow \ egin{array}{l} E'(t) \leq 0 \ \Rightarrow \ E(t) \leq E(0) \ \\ E'(t) \geq rac{2a'(t)}{a(t)} E(t) \ \\ \Rightarrow \ E(t) \geq \exp\left(2\lograc{a(t)}{a(0)}
ight) E(0) = rac{a(t)^2}{a(0)^2} E(0) \ \\ \left\{ a'(t) \leq 0 \ \Rightarrow \ E(0) \leq E(t) \leq rac{a_1^2}{a_0^2} E(0) \ a'(t) \geq 0 \ \Rightarrow \ rac{a_0^2}{a_1^2} E(0) \leq E(t) \leq E(0) \end{array}
ight.$$

Theorem 0.1. If a(t) is a monotone increasing or decreasing functions, then GEC is valid.

Oscillating coefficients

$$-rac{2|a'(t)|}{a(t)}E(t) \leq E'(t) \leq rac{2|a'(t)|}{a(t)}E(t)$$

$$\exp\left(-rac{2}{a_0}\int_0^t|a'(s)|ds
ight)E(0)\leq E(t)\leq \exp\left(rac{2}{a_0}\int_0^t|a'(s)|ds
ight)E(0)$$

Theorem 0.2. If $a(t) \in C^1([0,\infty))$ satisfies

$$|a'(t)| \le C(1+t)^{-\beta} \ (\beta \ge 0),$$

then (3) is valid for

$$\eta(t) = egin{cases} \exp\left(Ct^{-eta+1}
ight) & (eta < 1), \ t^C & (eta = 1), \ C & (eta > 1). \end{cases}$$

Periodic coefficients

Theorem 0.3. If a(t) is positive, periodic and non-constant, then for any $\varepsilon > 0$ the following uniform estimates are not valid in general:

$$E(t_1) \leq E(t_0) \exp\left(C\left(t_1^{1-arepsilon} - t_0^{1-arepsilon}
ight)
ight) \ \ 0 < orall t_0 < orall t_1 < \infty.$$

Example.
$$a(t) = 2 + \cos t$$

$$\limsup_{t \to \infty} \{|a'(t)|\} = 1 \implies e^{-Ct} E(0) \le E(t) \le e^{Ct} E(0)$$

Observations

• Can we derive some cancellation of the energy due to the oscillating coefficient?

$$a'(t) < 0 \Rightarrow E'(t) \le 0, \ a'(t) > 0 \Rightarrow E'(t) \ge 0.$$

• The order of error $\eta(t)$ is described by the number of oscillation:

$$\int_0^t |a'(s)| ds.$$

 \circ L^1 property of a'(t) concludes GEC; $\eta(t)=1$.

 $a'(t) \not\in L^1$

Theorem 1([5]). If $a(t) \in C^2([0,\infty))$ satisfies:

$$|a'(t)| \le C_1(1+t)^{-1}, \ |a''(t)| \le C_2(1+t)^{-2},$$

then GEC is valid.

Remark. $a(t) = 2 + \cos(\omega(t)), \, \omega(0) = 0, \, \omega'(t) > 0.$ $a' \in L^1 \Leftrightarrow \int_0^\infty \omega'(s) ds = \lim_{t \to \infty} \omega(t) < \infty$ \Leftrightarrow finite oscillation.

$$\omega(t) = \log(1+t) \Rightarrow egin{cases} |a^{(k)}(t)| \leq C_k (1+t)^{-k} \; (k=1,2); \ a'
ot\in L^1. \end{cases}$$

- GEC can be valid for infinitely oscillating coefficient.
- Cancellation of the energy is realized.

Stabilization property and C^m property

$$a \in C^m \ (m \geq 2), \ \lim_{t \to \infty} rac{1}{t} \int_0^t a(s) ds = \exists a_{\infty}.$$

We introduce the stabilization property with $\alpha \in [0, 1)$ and the C^m property with $\beta \in [0, 1)$ by

$$\int_{0}^{t} |a(s) - a_{\infty}| ds = O(t^{\alpha}) \quad (t \to \infty)$$
(Stabilization property)

$$|a^{(k)}(t)| \le C_k (1+t)^{-k\beta} \quad (k=1,\cdots,m)$$
 (6)
(C^m property)

Theorem 2([2]). If (5) and (6) are valid, then (3) holds for

$$\eta(t) = \exp\left(C(1+t)^{\sigma_m}\right),\,$$

$$\sigma_m = \max\left\{0,\, lpha - eta + rac{1-lpha}{m}
ight\}.$$

Remark.

- σ_m is monotone decreasing with respect to m and β , and monotone increasing with respect to α .
- GEC holds if $\beta \geq \alpha + \frac{1-\alpha}{m}$.
- The estimates (E) cannot be improved for $\beta < \alpha$.
- $\alpha = \beta$ is the critical case for GEC.

Orders of oscillating speed |a'(t)| and error $\eta(t)$

$$|a'(t)| \leq C(1+t)^{-eta}$$
 $\eta(t)^{-1}E(0) \leq E(t) \leq \eta(t)E(0)$

$$\int_0^t |a(s)-a_\infty|ds \leq C(1+t)^lpha, \; \eta(t)=\exp(C(1+t)^{\sigma_m})$$

Singular

Regularity of \boldsymbol{a}

Regular

m	1	2	• • •	m	• • •	∞
σ_m	1-eta	$rac{lpha+1}{2}-eta$	• • •	$\alpha - \beta + \frac{1-\alpha}{m}$	• • •	$\alpha - eta$

Unstable

Behavior of $\boldsymbol{E}(t)$

Stable

$$(0 \le \alpha \le \beta < 1)$$

2. Main results

We have the following questions from Theorem 2 as $m \to \infty$:

 $\cdot \alpha < \beta \Rightarrow \exists m \in \mathbb{N} \text{ s.t. GEC holds. Does GEC hold for } \alpha = \beta$?

$$\cdot \ lpha = eta \Rightarrow \eta(t) = \exp\left(C(1+t)^{rac{1-lpha}{m}}
ight)$$

What about the order of $\eta(t)$ in the limit case $m = \infty$?

We consider such problems to introduce the Gevrey class of a(t).

$$a(t) \in \gamma_{\rho}^{\nu} \Leftrightarrow |a^{(k)}(t)| \le C\rho^{-k}k!^{\nu} \quad (\nu > 1, \, \rho > 0)$$

$$C^{\omega} \subset \gamma_{\rho}^{\nu} \subset C^{\infty}$$

We consider the estimates (3) near the critical case $\alpha = \beta$ to introduce the following conditions:

$$\int_0^t |a(s) - a_{\infty}| ds = O(t^{\alpha}) \quad (t \to \infty)$$
 (5)

$$|a^{(k)}(t)| \le Ck!^{\nu} \left((1+t)^{\alpha} \left(\log(e+t) \right)^{\delta} \right)^{-k}$$

$$(k \in \mathbb{N}, \ \delta \ge 0)$$
(8)

$$egin{aligned} eta & eta & eta \ |a^{(k)}(t)| \leq C_k (1+t)^{-keta} \ \ (k \in \mathbb{N}, \ lpha < eta) \end{aligned}$$

<u>Theorem 3([3])</u>. If **(5)** and **(8)** are valid, then **(3)** holds for

$$\eta(t) = \exp\left(C\left(\log(e+t)\right)^{\sigma}\right), \ \ \sigma = \max\{0, \ \nu - \delta\}.$$

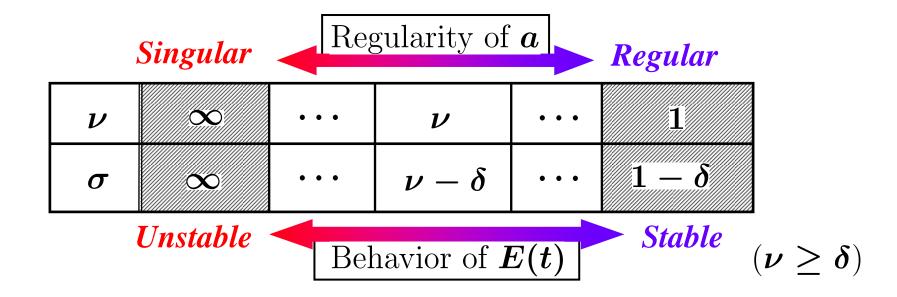
Theorem 3([3]). If (5) and (8) are valid, then (3) holds for

$$\eta(t) = \exp\left(C\left(\log(e+t)\right)^{\sigma}\right), \quad \sigma = \max\{0, \nu - \delta\}.$$

$$\eta(t)^{-1}E(0) \le E(t) \le \eta(t)E(0)$$
(3)

$$\int_0^t |a(s) - a_{\infty}| ds = O(t^{\alpha}) \quad (t \to \infty)$$
 (5)

$$|a^{(k)}(t)| \le Ck!^{\nu} \left((1+t)^{\alpha} \left(\log(e+t) \right)^{\delta} \right)^{-k}$$
 (8)



Summary

$$\int_0^t |a(s)-a_\infty|ds=O(t^lpha)$$
 $|a^{(k)}(t)| \leq C_k(1+t)^{-klpha} \; (k=1,\cdots,m)$ $\int_0^t |a(s)-a_\infty|ds=O(t^lpha)$

	m <	$< \infty$	$m=\infty$	$C_k = C k!^ u$			
a(t)	C^1	C^m	C^{∞}	$\gamma^{ u}$ C^{ω}			
$oxed{\eta(t)}$	$\exp(Ct^{1-lpha})$	$\exp(Ct^{rac{1-lpha}{m}})$	$\exp(Ct^arepsilon)$	$\exp(C(\log t)^ u)$ t^C			

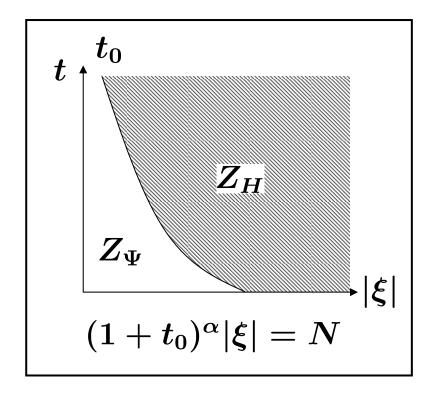
$$(0 \le \alpha < 1, \nu > 1)$$

3. Keys of the proof

- Refined diagonalization
- Division of infinitely many zones
- Algebra of the Gevrey functions for symbol calculus

$$(\partial_t^2 + a(t)^2 |\xi|^2) v(t,\xi) = 0$$





$$\partial_t V_1 = (\Phi_1 + B_1) V_1, \;\; V_1 = \left(egin{array}{c} v_1 \ \overline{v_1} \end{array}
ight) = \left(egin{array}{c} \partial_t v + ia |\xi| v \ \partial_t v - ia |\xi| v \end{array}
ight)$$

$$\Phi_1 = \left(egin{array}{cc} rac{a'}{2a} + ia|\xi| & 0 \ 0 & rac{a'}{2a} - ia|\xi| \end{array}
ight), \;\; B_1 = \left(egin{array}{cc} 0 & -rac{a'}{2a} \ -rac{a'}{2a} & 0 \end{array}
ight)$$

$$egin{aligned} egin{aligned} egin{aligned} \partial_t V_1 &= (\Phi_1 + B_1) V_1 \ \end{pmatrix} & \Phi_1 &= \left(egin{array}{cc} \phi_1 & 0 \ 0 & \overline{\phi_1} \end{array}
ight), \;\; B_1 &= \left(egin{array}{cc} 0 & \overline{b_1} \ b_1 & 0 \end{array}
ight) \ V_2 &= M_1^{-1} V_1, \;\; M_1 &= \left(egin{array}{cc} 1 & \overline{\delta_1} \ \delta_1 & 1 \end{array}
ight), \;\; \delta_1 &= rac{-i b_1}{2 \phi_{1, \Im}} \end{aligned}$$

$$|\delta_1| = rac{|rac{a'}{2a}|}{2a|\xi|} \leq rac{C_1(1+t)^{lpha-eta}}{4a_0^2N} \leq rac{C_1}{4a_0^2N} \leq rac{1}{2} \ (lpha \leq eta, \;\; N \gg 1)$$

$$oxed{\partial_t V_2 = (\Phi_2 + B_2) V_2} \quad \Phi_2 = \left(egin{array}{cc} \phi_2 & 0 \ 0 & \overline{\phi_2} \end{array}
ight), \ B_2 = \left(egin{array}{cc} 0 & \overline{b_2} \ b_2 & 0 \end{array}
ight)$$

$$\phi_{2,\Re} = rac{1}{2} \partial_t \left(\log \left(rac{a}{1 - |\delta_1|^2}
ight)
ight), \;\; \phi_{2,\Im} = a |\xi| - rac{2 |\delta_1|^2}{1 - |\delta_1|^2}$$

$$\partial_t V_j = (\Phi_j + B_j) V_j$$



$$V_j = \left(egin{array}{c} v_j \ \overline{v_j} \end{array}
ight), \Phi_j = \left(egin{array}{c} \phi_j & 0 \ 0 & \overline{\phi_j} \end{array}
ight), \ B_j = \left(egin{array}{c} 0 & \overline{b_j} \ b_j & 0 \end{array}
ight)$$

$$V_{j+1}=M_j^{-1}V_j,\;\;M_j=\left(egin{array}{cc}1&\overline{\delta_j}\\delta_j&1\end{array}
ight),\;\;\delta_j=rac{-ib_j}{2\phi_{j,\Im}}$$

$$\partial_t V_{j+1} = (\Phi_{j+1} + B_{j+1}) V_{j+1}$$

$$\begin{cases} \phi_{j+1,\Re} = \frac{1}{2} \partial_t \left(\log \left(\frac{a}{\prod_{k=1}^j (1 - |\delta_k|^2)} \right) \right) \\ \phi_{j+1,\Im} = a |\xi| + \sum_{k=1}^j \frac{-2|\delta_k|^2 \phi_{k,\Im} + \Im\{\delta'_k \overline{\delta_k}\}}{1 - |\delta_k|^2} \\ b_{j+1} = \frac{b_j |\delta_j|^2 - \delta'_j}{1 - |\delta_j|^2} \qquad (j = 0, \dots, m - 1) \end{cases}$$

$$|V_m(t,\xi)|^2 \left\{egin{array}{l} \leq |V_m(t_0,\xi)|^2 \exp\left(2\int_{t_0}^t \left(\phi_{m,\Re} + |b_m|
ight) \; ds
ight) \ \geq |V_m(t_0,\xi)|^2 \exp\left(2\int_{t_0}^t \left(\phi_{m,\Re} - |b_m|
ight) \; ds
ight) \end{array}
ight.$$

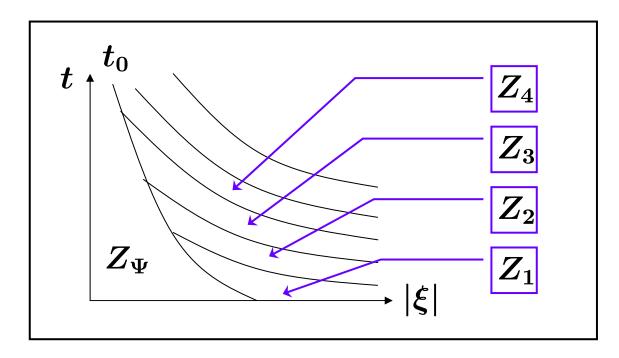
$$|V_m(t,\xi)|^2 \simeq |V_1(t_0,\xi)|^2 \ (|M_k-I| \ll 1, \ k=1,2,\cdots,m)$$
 $|b_k(t,\xi)| \leq C_k |\xi|^{-k+1} (1+t)^{-\beta k} \ (k=1,2,\cdots,m)$

$$\int_{0}^{t} |\xi|^{-k+1} (1+s)^{-\beta k} ds = \frac{1}{\beta k - 1} |\xi|^{-k+1} (1+t_0)^{-\beta k + 1}$$
$$= \frac{N^{-k+1}}{\beta k - 1} (1+t_0)^{-k(\beta - \alpha) + 1 - \alpha}$$

The estimates of $|V_1|$ is improved by diagonalization procedures due to M_k .

Division of infinitely many zones

$$Z_k = \{(t, \xi) \; ; \; t_{k-1} \le t \le t_k \} \; ,$$
 $(1 + t_k)^{lpha} \left(\log(e + t_k) \right)^{\delta} |\xi| = (k+1)^{
u}$



Algebra of the Gevrey functions

$$\sum_{k=0}^{n} \binom{n}{k} \left(\frac{k!(n-k)!}{n!} \right)^{\nu} \le C \quad \Leftrightarrow \quad \nu > 1$$